

# Asymptotics

## Big O notation

$$f(x) = O(g(x)) \text{ as } x \rightarrow a \Leftrightarrow \left| \frac{f(x)}{g(x)} \right| \text{ bdd as } x \rightarrow a$$

## little o notation

$$f(x) = o(g(x)) \text{ as } x \rightarrow a \Leftrightarrow \left| \frac{f(x)}{g(x)} \right| \rightarrow 0 \text{ as } x \rightarrow a$$

## Capital O

$$f(x) = O(g(x)) \text{ as } x \rightarrow a \Leftrightarrow f = O(g) \text{ \& } g = O(f)$$

## Asymptotic definitions

Series:  $f(x) \sim \sum_{j=1}^N f_j(x) \text{ as } x \rightarrow a$

$$\forall \text{ finite } m \in \mathbb{N}, f(x) - \sum_{j=1}^m f_j(x) = o(f_m(x))$$

[  $\varepsilon^2 = o(\varepsilon^2)$  ]

Sequence:  $(a_n)_{n \geq 1}$  s.t.  $a_{n+1}(x) = o(a_n(x)) \text{ as } x \rightarrow a$

Given a sequence  $(a_n)_{n \geq 1}$ , the asymptotic series

$$f \sim \sum_{j=1}^N f_j a_j(x), \text{ then } f_j \text{ coeff. unique}$$

$$f(x, \varepsilon) \sim \sum_{j=1}^N f_j(x) a_j(\varepsilon) \text{ as } \varepsilon \rightarrow 0 \text{ is a Poincaré expansion}$$

Poincaré expansions have unique coefficients  $f_j$ , given  $(a_n)_{n \geq 1}$

## Singular Expansions

(all pretty general, particular to integrals)

- Rescale 1st. (include terms left out when set  $\varepsilon=0$ )
- For non-linear problems, rescale both  $f$  &  $x$
- For some problems, need different rescalings in different places.
  - use correct scaling in correct place & match

## Matching

## Matching M

If we have different expansions in different places, need to match:

(a) Ad hoc: look at it by eye (few checks could be wrong)

(b) Van Dyke's Matching Rule (Sheet 4 Q5)  
 $E_P H_Q f = H_Q E_P f$

(c) intermediate variables (did this w/ integrals (Sheet 2 Q2,3), but non-exam for ODEs).

Composite expansions:  $C = E_P f + H_Q f - E_P H_Q f$   
 (not Poincaré  $\Rightarrow$  not unique, Sheet 4 Q5)

## Endpoints:

Watson's Lemma

$$I(\lambda) \sim \frac{1}{\lambda} f(A) e^{-\lambda g(A)} \text{ (simple cases)}$$

## interior:

Laplace's method:  $g'(x_0) = 0$

$$I(\lambda) \sim \sqrt{\frac{2\pi}{-\lambda g''(x_0)}} f(x_0) e^{\lambda g(x_0)}$$

## Regular Expansions

(1) Iteration method

(2) Successive approximation

- set  $\varepsilon=0$ , solve for  $f_0$
- set  $f = f_1 + \varepsilon_1$ , found exact eqn for  $\varepsilon_1$
- Approx solve for  $\varepsilon_1$  (by setting  $\varepsilon=0$ ), to get  $\hat{\varepsilon}_1$ .  $f = f_0 + \hat{\varepsilon}_1 + \varepsilon_2, \dots$

(3) Series Expansion

- Guess series  $f = a_0(\varepsilon) f_0 + a_1(\varepsilon) f_1 + \dots$
  - Substitut in & solve for each  $f_0, f_1, f_2, \dots$
- e.g.  $f = f_0 + \varepsilon f_1 + \varepsilon^2 f_2 + \dots$

If it's a regular problem. Do the above & it will work. Problem for Singular situations.

If you set  $\varepsilon=0$ , lose leading order terms.

$\Rightarrow$  so RESCALE to keep those important terms

## Choice of rescaling

- Balance terms in equations or integrals (Sheet 1 Q4)
- Solve for  $\varepsilon=0$  & look for where terms are larger or smaller than expected (Sheet 4 Q3 & Q4)
- Solve leading & 1st order & see where the correction term becomes large.

E.g.  $f \sim 1 + \varepsilon x$

$x = O(1) \rightarrow \checkmark$

$x = O(\frac{1}{\varepsilon}) \rightarrow$  correction term too big.

General Asymptotics  $\Rightarrow$  Specialise

## Integrals w/ exponentials

$$I(\lambda) = \int_A^B f(x) e^{\lambda g(x)} dx \text{ as } \lambda \rightarrow \infty$$

w/ integral, more information than you need. Different integrands give the same number (the integral - more degrees of freedom).

look at where integrand exponentially decays

Steepest descent:

deform contour s.t.  $g'(z_*) = 0$  (saddle pt) Then use Laplace's method

## Stationary phase

$$I(\lambda) = \int_A^B f(x) e^{i\lambda g(x)} dx \text{ highly oscillatory}$$

$\Rightarrow$  set  $i\tilde{g}(x) = g(x)$  & use steepest descent

# Integral Transforms

extended  $f: [0, \infty) \rightarrow \mathbb{C}$  to  $g: \mathbb{R} \rightarrow \mathbb{C}$   
by zero extending  $g(t) = \begin{cases} f(t) & t \geq 0 \\ 0 & t < 0 \end{cases}$

## Fourier Transforms

'forward'  $\tilde{f}(\omega) = \int_{-\infty}^{\infty} f(t) e^{-i\omega t} dt$

inverse  $f(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \tilde{f}(\omega) e^{i\omega t} d\omega$

## Fourier Properties

name	$g(t)$	$\tilde{g}(\omega)$
Shift 1	$f(t-a)$	$e^{-i\omega a} \tilde{f}(\omega)$
Shift 2	$e^{i\omega t} f(t)$	$\tilde{f}(\omega - \omega)$
Scaling	$f(at)$	$\frac{1}{ a } \tilde{f}\left(\frac{\omega}{a}\right)$
Diff	$f'(t)$	$i\omega \tilde{f}(\omega)$
$\times t$	$t f(t)$	$i \frac{d\tilde{f}}{d\omega}(\omega)$
convolution	$\int_{-\infty}^{\infty} f(\tau) g(t-\tau) d\tau$	$\tilde{f}(\omega) \tilde{g}(\omega)$

## Plancherel's Theorem

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} |\tilde{f}(\omega)|^2 d\omega = \int_{-\infty}^{\infty} |f(t)|^2 dt$$

energy in frequency density      energy in time signal

## Laplace Transforms

$$\hat{f}(s) = \int_0^{\infty} f(t) e^{-st} dt$$

$$f(t) = \frac{1}{2\pi i} \int_{-i\infty+a}^{i\infty+a} \hat{f}(s) e^{st} ds$$

$\times$  pass all singularities

## Laplace Properties

name	$g(t)$	$\hat{g}(s)$
Shift 1	$f(t-a)$	$e^{-sa} \hat{f}(s)$
Shift 2	$e^{at} f(t)$	$\hat{f}(s-a)$
Scaling	$f(at)$	$\frac{1}{a} \hat{f}\left(\frac{s}{a}\right)$
Diff	$f'(t)$	$-\hat{f}(0) + s \hat{f}(s)$
	$f''(t)$	
$\times t$	$t f(t)$	$-\frac{d\hat{f}}{ds}(s)$
	$t^n$	$\frac{n!}{s^{n+1}}$
convolution	$\int_0^{\infty} f(\tau) g(t-\tau) d\tau$	$\hat{f}(s) \hat{g}(s)$

## Examples

$f(t)$	$\hat{f}(s)$
1	$\frac{1}{s}$
$e^{at}$	$\frac{1}{s-a}$
$\cos(at)$	$\frac{s}{s^2+a^2}$
$\sin(at)$	$\frac{a}{s^2+a^2}$
$t^n$	$\frac{n!}{s^{n+1}}$

$$\frac{d\hat{y}}{dt} = -y(0) + s \hat{y}(s) \quad \text{JUST OIFF}$$

$$\begin{aligned} \frac{d^2\hat{y}}{dt^2} &= -\dot{y}(0) + s \frac{d\hat{y}}{ds} \\ &= -\dot{y}(0) + s(-y(0) + s \hat{y}(s)) \\ &= -\dot{y}(0) - sy(0) + s^2 \hat{y}(s) \end{aligned}$$